

q -oscillator lattice and representations of quantum groups

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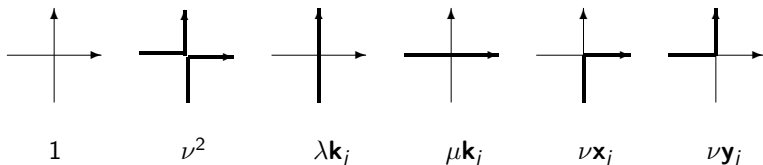
New frontiers in exactly solved models, ANU, July 21, 2005



- 1 Formulation of q -oscillator lattice
- 2 Tetrahedron equation
- 3 Representations of quantum groups

Weights

Let \mathbf{Z} be a partition function of completely inhomogeneous ice-type model on a rectangular periodical lattice $N \times M$ with the following weights of j^{th} vertex:



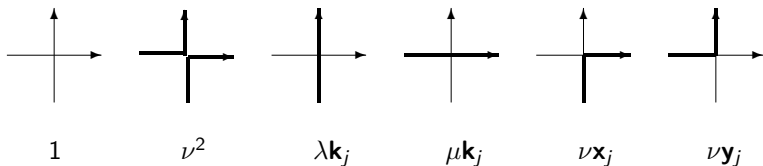
Numbers λ, μ [$\nu^2 \equiv -q^{-1}\lambda\mu$] are the same for all vertices.

$$\mathcal{H} : \quad \mathbf{x}\mathbf{y} = 1 - q^2\mathbf{k}^2, \quad \mathbf{y}\mathbf{x} = 1 - \mathbf{k}^2, \quad \mathbf{x}\mathbf{k} = q\mathbf{k}\mathbf{x}, \quad \mathbf{k}\mathbf{y} = q\mathbf{y}\mathbf{k}.$$

Vertex index j stands for a component of tensor power $\mathbf{k}_j, \mathbf{x}_j, \mathbf{y}_j \in \mathcal{H}^{\otimes NM}$.

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$$\mathcal{H} : \mathbf{xy} = 1 - q^2 \mathbf{k}^2, \quad \mathbf{yx} = 1 - \mathbf{k}^2, \quad \mathbf{xk} = q \mathbf{kx}, \quad \mathbf{ky} = q \mathbf{yk}.$$

Vertex index j stands for a component of tensor power $\mathbf{k}_j, \mathbf{x}_j, \mathbf{y}_j \in \mathcal{H}^{\otimes NM}$.

Partition function is a polynomial of λ, μ with operator-valued coefficients

$$\mathbf{Z}(\lambda, \mu) = \sum_{n,m=0}^{M,N} \lambda^{Mn} \mu^{Nm} \mathbf{z}_{n,m}, \quad \mathbf{z}_{n,m} \in \mathcal{H}^{\otimes NM}$$

Theorem

$$\mathbf{Z}(\lambda, \mu) \mathbf{Z}(\lambda', \mu') = \mathbf{Z}(\lambda', \mu') \mathbf{Z}(\lambda, \mu)$$

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Matrix L

$$L_{\alpha,\beta}(\mathcal{H}; \lambda, \mu) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \lambda \mathbf{k} & \nu \mathbf{y} & 0 \\ 0 & \nu \mathbf{x} & \mu \mathbf{k} & 0 \\ 0 & 0 & 0 & \nu^2 \end{pmatrix}, \quad \nu^2 = -q^{-1} \lambda \mu$$

$L_{\alpha,\beta}$ acts in $V_\alpha \otimes V_\beta \otimes F$, $V = \mathbb{C}^2$ and F is e.g. the Fock space $\mathbf{k} = q^n$.

$$L_{\alpha,\beta} = \prod_n \prod_m L_{\alpha_n, \beta_m}(\mathcal{H}_{n,m}; \lambda, \mu), \quad \mathbf{Z} = \text{Trace}_{V_\alpha, V_\beta}(L_{\alpha,\beta})$$

$$\prod_n f_n = f_N f_{N-1} \dots f_1, \quad V_\alpha = \bigotimes_{n=1}^N V_{\alpha_n} \text{ etc.}$$

Tetrahedron equation

$$M_{\alpha,\alpha'} \left(\mathcal{H}_0; \frac{\mu}{\mu'} \right) M_{\beta,\beta'} \left(\mathcal{H}_0; \frac{\lambda'}{\lambda} \right) L_{\alpha,\beta}(\mathcal{H}; \lambda, \mu) L_{\alpha',\beta'}(\mathcal{H}; \lambda', \mu') = \\ L_{\alpha',\beta'}(\mathcal{H}; \lambda', \mu') L_{\alpha,\alpha'}(\mathcal{H}; \lambda, \mu) M_{\beta,\beta'} \left(\mathcal{H}_0; \frac{\lambda'}{\lambda} \right) M_{\alpha,\alpha'} \left(\mathcal{H}_0; \frac{\mu}{\mu'} \right)$$

where

$$M(\mathcal{H}_0; \xi) = \begin{pmatrix} \xi^{\mathbf{n}_0} & 0 & 0 & 0 \\ 0 & \lambda_0(-q\xi)^{\mathbf{n}_0} & \nu_0 \xi^{-1/2+\mathbf{n}_0} \mathbf{y}_0 & 0 \\ 0 & \nu_0 \xi^{1/2+\mathbf{n}_0} \mathbf{x}_0 & \mu_0(-q\xi)^{\mathbf{n}_0} & 0 \\ 0 & 0 & 0 & \nu_0^2 \xi^{\mathbf{n}_0} \end{pmatrix}$$

$$\mathbf{k}_0 = q^{\mathbf{n}_0}, \nu_0^2 \equiv q^{-1} \lambda_0 \mu_0.$$

Lax operators and R matrices

$$\mathcal{L}_{\beta}^{(n)}(\lambda, \mu) = \text{Trace}_{V_{\alpha}} \prod_m^{\curvearrowright} L_{\alpha_n, \beta_m}(\mathcal{H}_{n,m}; \lambda, \mu) : V^{\otimes M} \otimes F^{\otimes M}$$

$$V^{\otimes M} = \bigoplus_{m=0}^M V_{\pi_m} \Rightarrow \mathcal{L}_{\beta}(\lambda, \mu) = \bigoplus_{m=0}^M \mu^m L_{\pi_m}(\lambda)$$

$\pi_m - m^{\text{th}}$ fundamental representation of $\mathcal{U}_q(\widehat{\mathfrak{sl}}_M)$, π_0 and $\pi_M -$ scalars. In quantum space $q^{\mathbf{n}_{n,1} + \mathbf{n}_{n,2} + \dots + \mathbf{n}_{n,M}} = q^J$ is the center, $F^{\otimes M} = \bigoplus_{J=0}^{\infty} V_{J\pi_1}$.

$$R_{\beta, \beta'} \left(\frac{\lambda}{\lambda'} \right) = \text{Trace}_{F_0} \prod_m^{\curvearrowright} M_{\beta_m, \beta'_m} \left(\mathcal{H}_0; \frac{\lambda'}{\lambda} \right) = \sum_{m, m'=0}^M \lambda_0^m \mu_0^{m'} R_{\pi_m, \pi_{m'}} \left(\frac{\lambda}{\lambda'} \right)$$

Lax operator for $M = 2$

$$\mathcal{L}(\lambda, \mu) = \begin{pmatrix} 1 + \lambda^2 q^J & & \\ & \nu^2 L_{1/2}(\lambda) & \\ & & \mu^2 (q^J + q^{-2} \lambda^2) \end{pmatrix}$$

where $q^J = q^{n_1 + n_2}$ is the center and

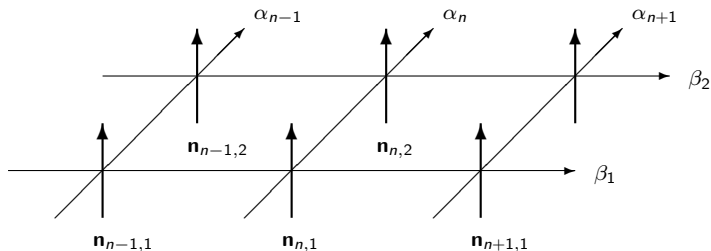
$$L_{1/2}(\lambda) = \begin{pmatrix} \lambda q^{n_2} - \lambda^{-1} q^{n_1+1} & \mathbf{x}_1 \mathbf{y}_2 \\ \mathbf{y}_1 \mathbf{x}_2 & \lambda q^{n_1} - \lambda^{-1} q^{n_2+1} \end{pmatrix}$$

In the limit $q \rightarrow 1$: $\mathbf{x} \rightarrow \mathbf{a}$, $\mathbf{y} \rightarrow \mathbf{a}^+$ and

$$J = \mathbf{a}_1^+ \mathbf{a}_1 + \mathbf{a}_2^+ \mathbf{a}_2, \quad H = \mathbf{a}_1^+ \mathbf{a}_1 - \mathbf{a}_2^+ \mathbf{a}_2, \quad E = \mathbf{a}_1^+ \mathbf{a}_2, \quad F = \mathbf{a}_2^+ \mathbf{a}_1$$

Six vertex model

$M = 2$, N – length of $\mathcal{U}_q(\widehat{sl}_2)$ chain.



Integrals of motion: $J_n = \mathbf{n}_{n,1} + \mathbf{n}_{n,2}$, $K_1 = \sum_n \mathbf{n}_{n,1}$, $K_2 = \sum_n \mathbf{n}_{n,2}$.

Six vertex model: $J_n = 1$ for all n (spin 1/2)

$\Rightarrow K_1 =$ “number of spins up”, $K_2 =$ “number of spins down”.

THANK YOU



Universal form of nested Bethe Ansatz

Decomposition of \mathbf{Z} :

$$\mathbf{Z}(\lambda, \mu) = \sum_{n,m=0}^{M,N} \lambda^{Mn} \mu^{Nm} \mathbf{z}_{n,m}, \quad \mathbf{z}_{n,m} \in \mathcal{H}^{\otimes NM}$$

Let $\mathbf{u}\mathbf{v} = q^2\mathbf{v}\mathbf{u}$ – additional auxiliary Weyl algebra and let

$$\mathbf{J}(\mathbf{u}, \mathbf{v}) = \sum_{n,m} (-q)^{-nm} \mathbf{u}^n \mathbf{v}^m \mathbf{z}_{n,m}$$

Universal NBA equation:

$$\langle Q | \mathbf{J}(\mathbf{u}, \mathbf{v}) = 0$$

E.g. $Q(u) = \langle Q | u \rangle$ etc.



Baxter's equation for $\mathcal{U}_q(\widehat{sl}_2)$

$$\mathbf{J}(\mathbf{u}, \mathbf{v}) = \sum_{n,m} (-q)^{-nm} \mathbf{u}^n \mathbf{v}^m z_{n,m} : \langle Q | \mathbf{J}(\mathbf{u}, \mathbf{v}) = 0$$

$$\text{For } M = 2 \quad J(\mathbf{u}, \mathbf{v}) = \phi(q^2 \mathbf{u}) - t(\mathbf{u}) \mathbf{v} + \phi'(q^{-2} \mathbf{u}) \mathbf{v}^2$$

$$\langle Q | u \rangle = Q(u), \quad \langle Q | \mathbf{u} | u \rangle = u Q(u), \quad \langle Q | \mathbf{v} | u \rangle = v Q(q^2 u)$$

$$\langle Q | J(\mathbf{u}, \mathbf{v}) | q^{-2} u \rangle \equiv \phi(u) Q(q^{-2} u) - t(u) Q(u) v + \phi'(u) Q(q^2 u) v^2 = 0$$

Fix v by condition $Q(0) = 1$, then $Q(u)$ is a polynomial.